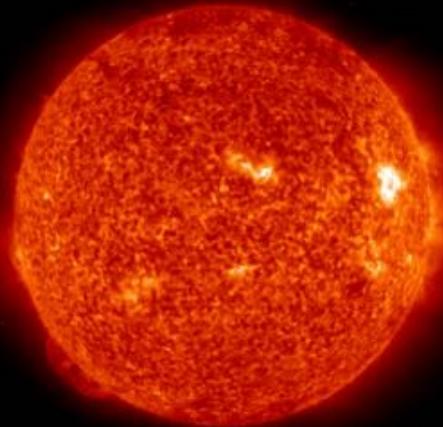
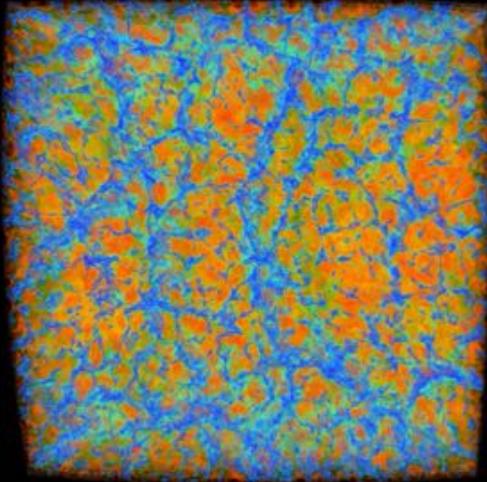


Simulations of Flux Emergence in Cool Stars: What's Convection, Rotation, and Stellar Structure got to do with it?

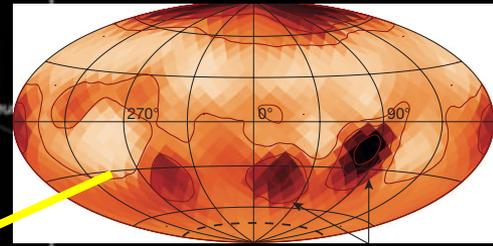
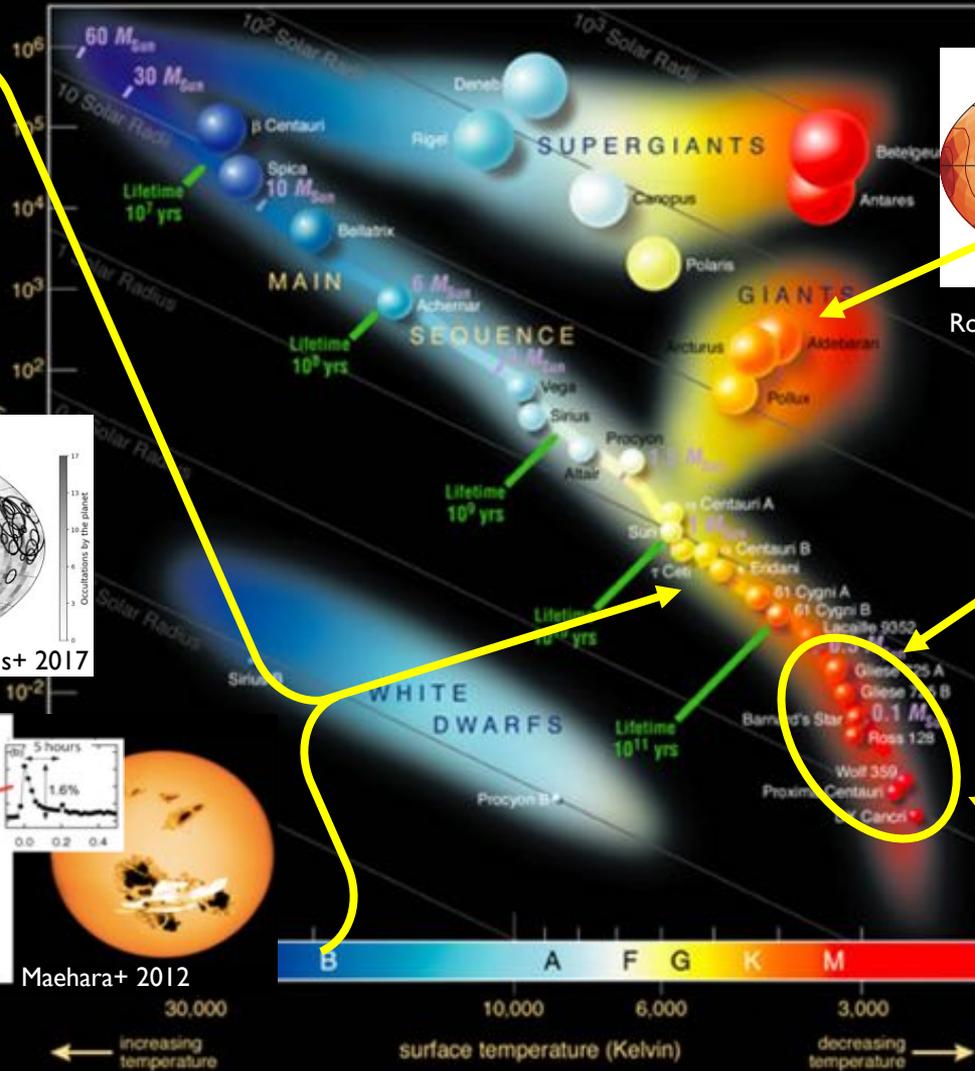


Maria Weber^{1,2}, **Matthew Browning**³, **Nicholas Nelson**⁵, **Yuhong Fan**⁶, **Mark Miesch**⁷, **Ben Brown**^{8,9},
⁴Suzannah Boardman, ⁴Joshua Clarke, ⁴Samuel Pugsley, ⁴Edward Townsend

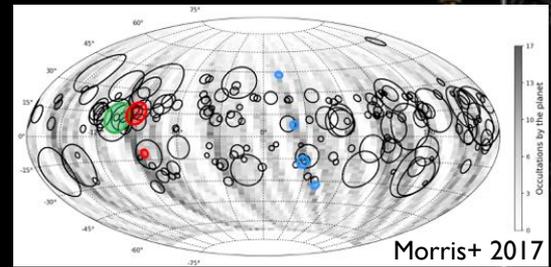
¹University of Chicago, ²Adler Planetarium, ³University of Exeter, ⁴former Mphys students at the University of Exeter,

⁵California State University, Chico, ⁶HAO/NCAR, ⁷NOAA, ⁸LASP,

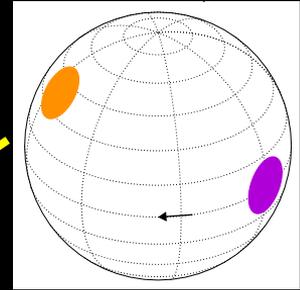
⁹Department of Astrophysical and Planetary Sciences, University of Colorado, Boulder



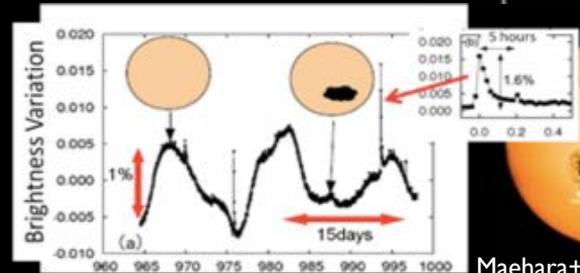
Roettenbacher+ 2016



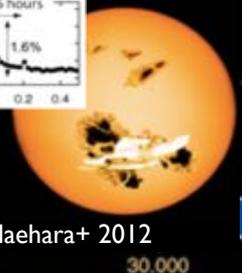
Morris+ 2017



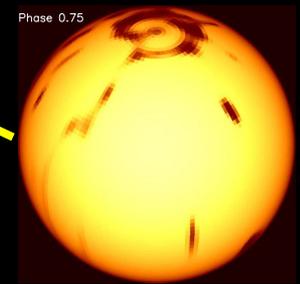
Davenport+ 2015



Maehara+ 2012



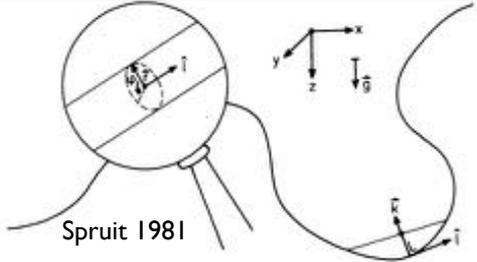
30,000



Barnes+ 2015

Model Schematic:

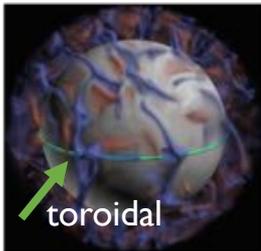
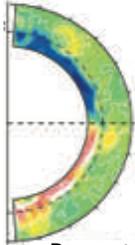
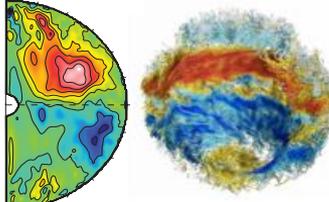


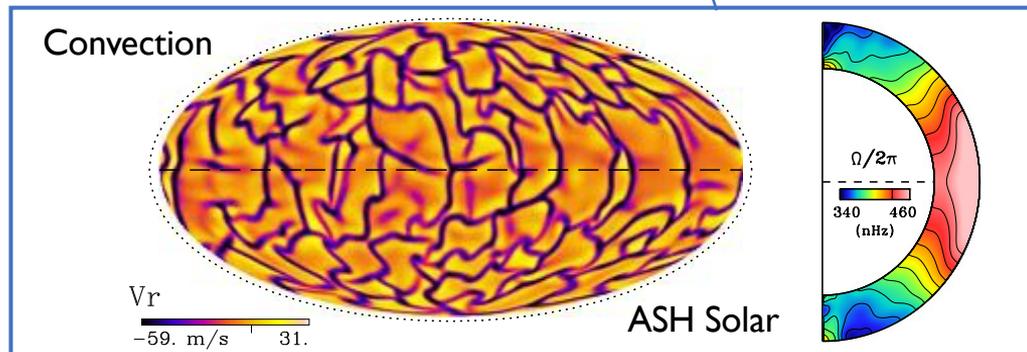
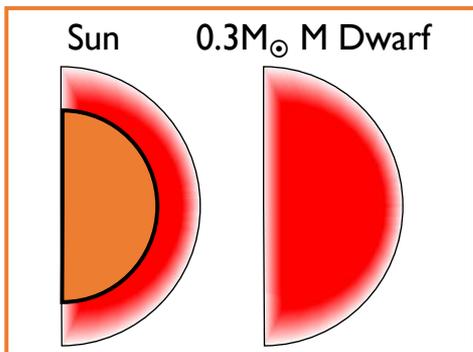


Equation of Motion:

$$\rho \frac{dv}{dt} = -2\rho(\boldsymbol{\Omega}_0 \times \mathbf{v}) - (\rho_e - \rho)\mathbf{g} + \mathbf{l} \frac{\partial}{\partial s} \left(\frac{B^2}{8\pi} \right) + \frac{B^2}{4\pi} \mathbf{k}$$

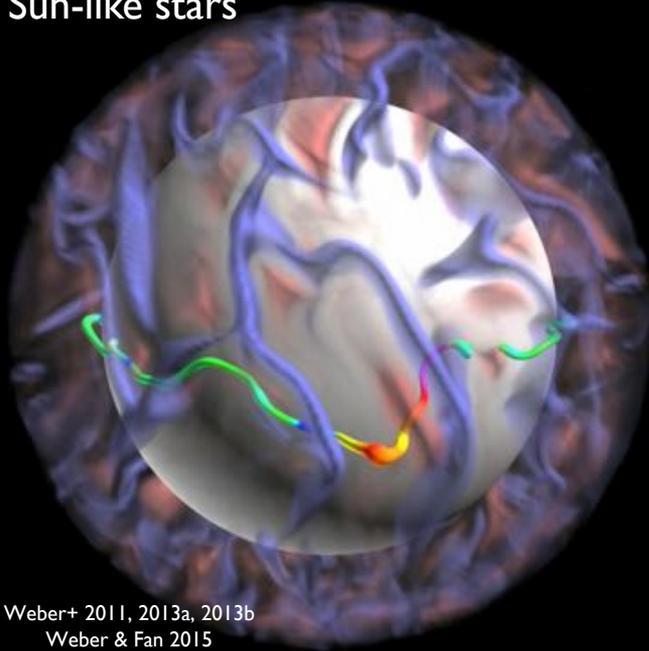
$$- C_d \frac{\rho_e |(\mathbf{v} - \mathbf{v}_e)_\perp| (\mathbf{v} - \mathbf{v}_e)_\perp}{(\pi \Phi / B)^{1/2}}$$

<p>B_0 r_0 θ_0 Φ</p> 	<p>Solar $1\Omega_\odot, 5\Omega_\odot$</p> <p>Tachocline interface dynamo</p>  <p>Browning+ 2006</p>	<p>M dwarf, Solar $3\Omega_\odot$</p> <p>Distributed dynamo</p>  <p>Browning 2008 M Dwarf Rapid Ω_0 Brown+ 2011</p>
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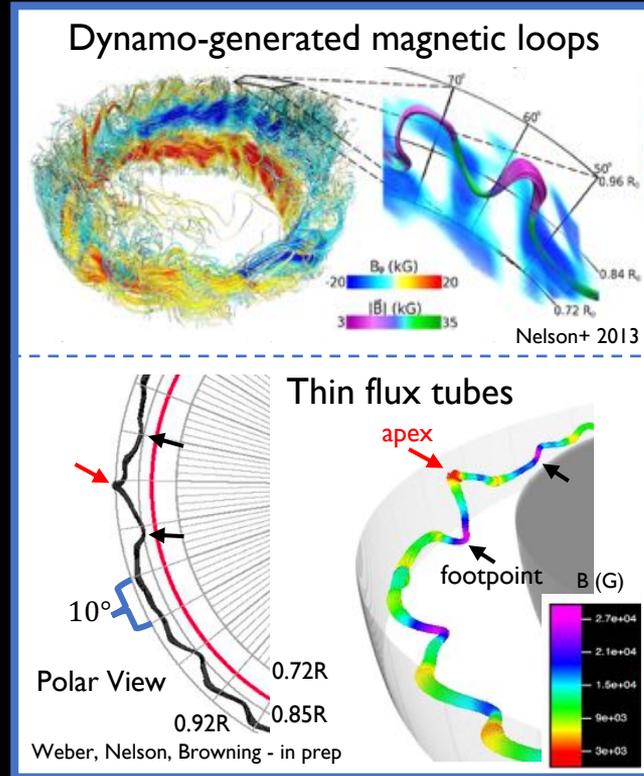
Convection modulates flux emergence

Sun-like stars

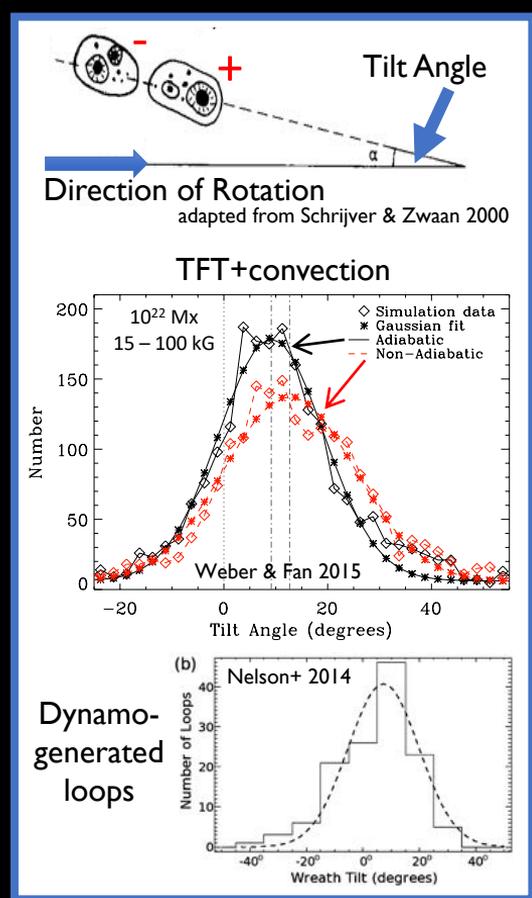


Weber+ 2011, 2013a, 2013b
Weber & Fan 2015

- Convection & magnetic buoyancy work in concert to promote flux emergence



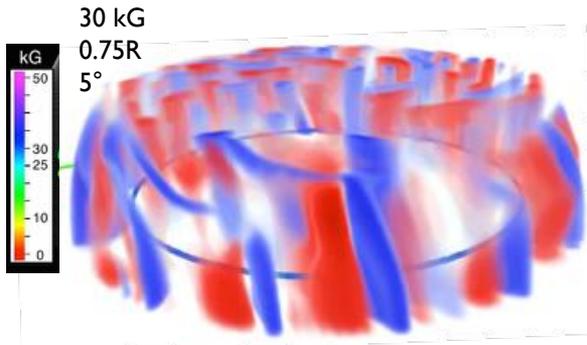
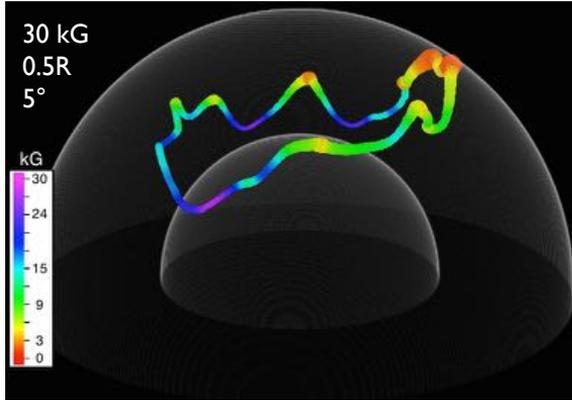
- Downflows naturally induce loops $\sim 15^\circ - 20^\circ$ apart



- Convection introduces a statistical spread in tilt angles

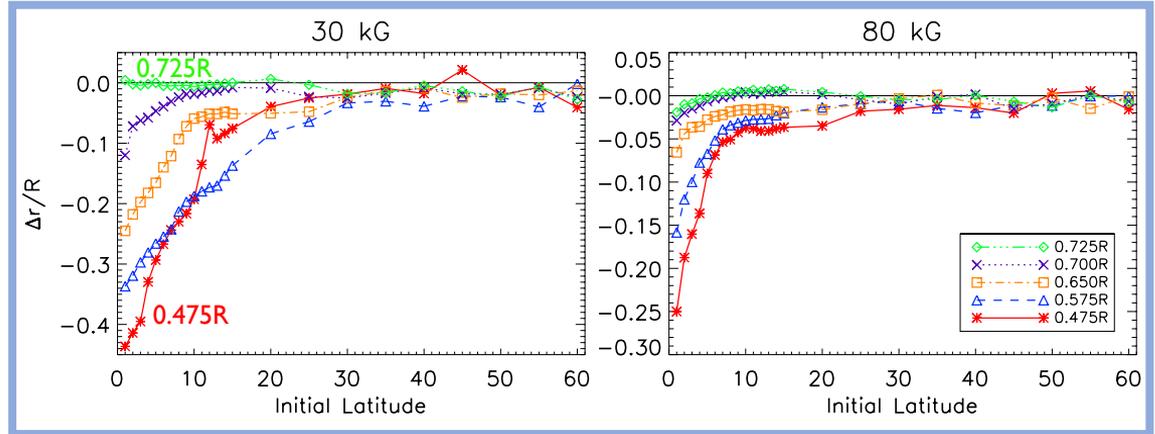
Convection can also suppress flux emergence

Fully convective M dwarf



Average 'suppression depth' Δr

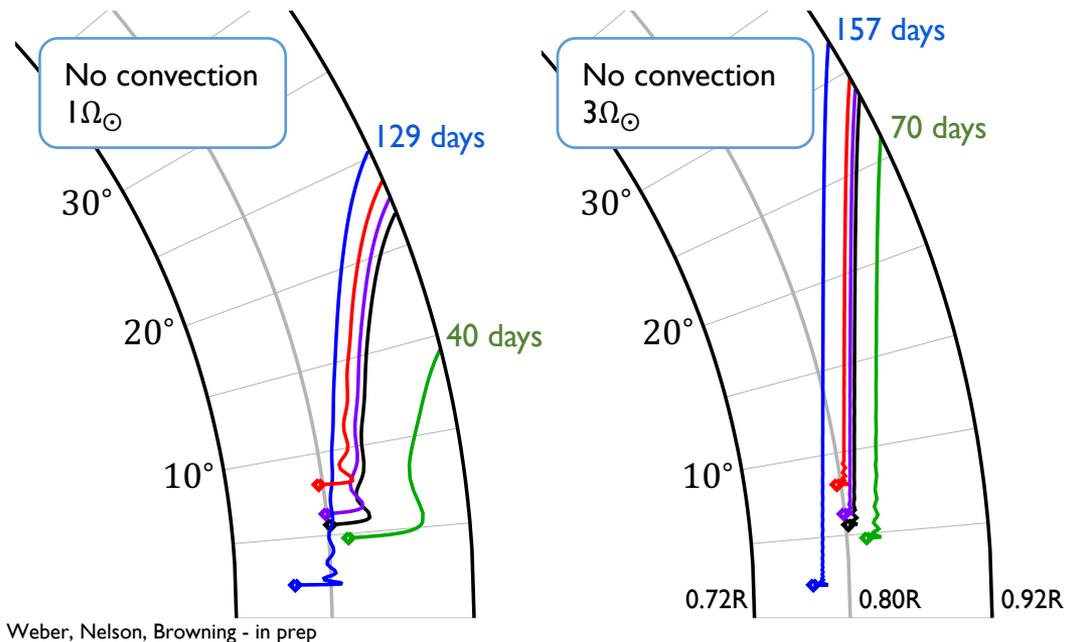
Weber+ 2017



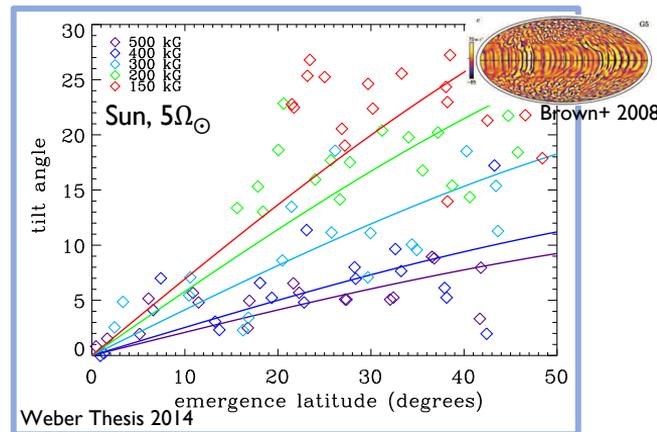
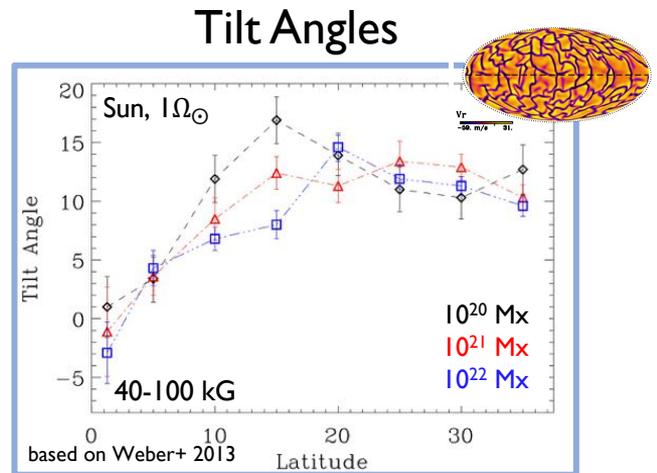
- The global rise of TFTs is more strongly suppressed by convective flows when the flux tube is initiated:
 - in the deeper interior
 - at lower latitudes
 - with a weaker magnetic field strength

Weber & Browning 2016, Weber+ 2017

Rotation alters emergence properties

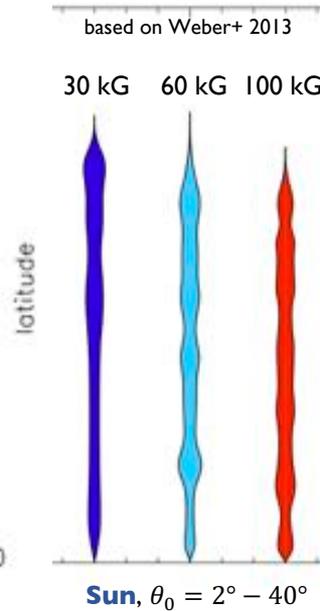
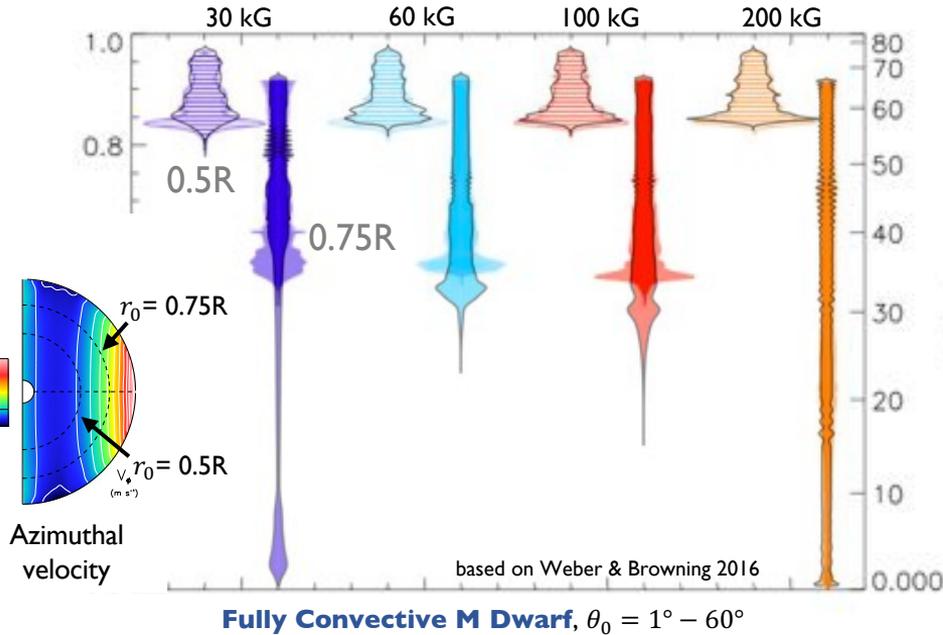


- Due to the Coriolis force, more rapid rotation:
 - Lengthens the rise time
 - Leads to poleward deflection
 - Increases tilt angles

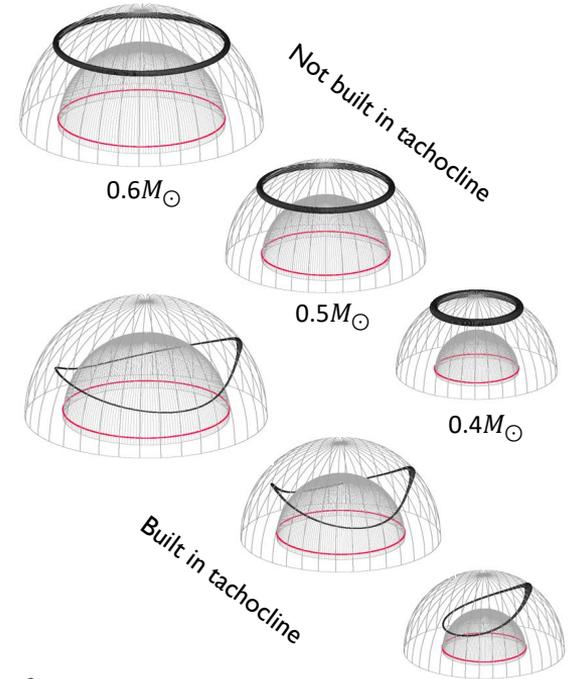


Stellar structure impacts emergence latitudes and more

Emergence Latitude Probability Functions



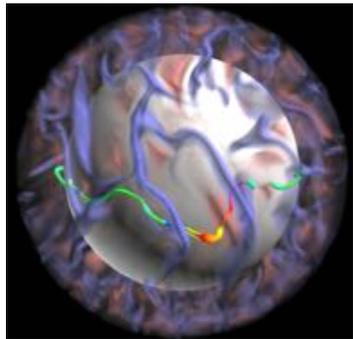
Tachocline or not?



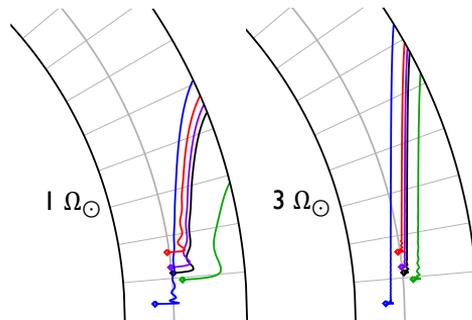
- Unlike solar case, in M dwarf there is a tendency for high latitude emergence ($> 30^\circ$)
- Exceptions when flux tubes initiated closer to the surface and of sufficiently weak (≤ 30 kG) or strong field strengths (≥ 200 kG)
- Increased density in M dwarfs leads to longer flux tube rise times by $\leq 10x$
- Assumption of flux tube generating region, and thereby initial thermodynamic properties, matter

Summary

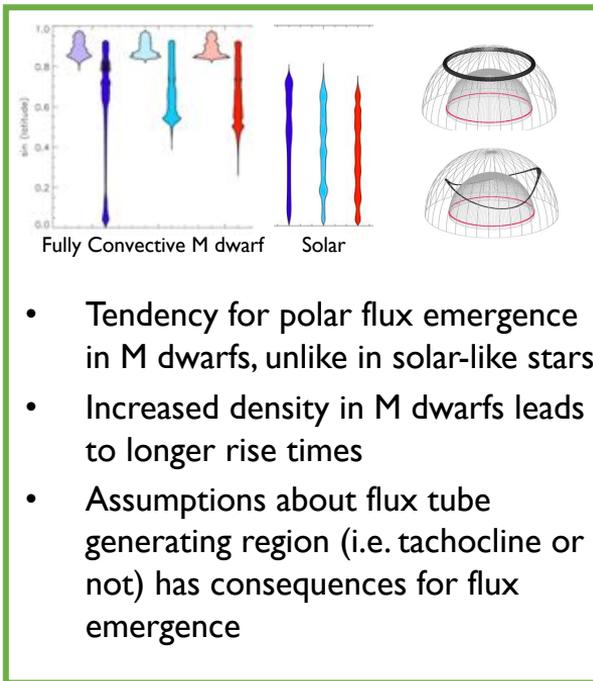
Convection, rotation, and stellar structure are all important contributing factors to the overall trend of flux emergence.



- Convection modulates flux emergence
- Fluid motions both suppress and promote the rise of magnetism
- Convection introduces a statistical spread in emergence properties



- Due to the Coriolis force, rapid rotation:
 - Lengthens the rise times
 - Leads to poleward emergence
 - Increases tilt angle



- Tendency for polar flux emergence in M dwarfs, unlike in solar-like stars
- Increased density in M dwarfs leads to longer rise times
- Assumptions about flux tube generating region (i.e. tachocline or not) has consequences for flux emergence

This work is a step toward linking magnetic flux emergence, convection, and dynamo action along the lower end of the main sequence.